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21 separate items enclosed

MEMORANDUM FOR PR (In-House Publication)

FROM: PROI (TI) (STINFO)

30 November 1999

SUBJECT: Authorization for Release of Technical Information, Control Number: **AFRL-PR-ED-TP-1999-0226**Talley, D., "Basic Research in Supercritical Combusion" (BFI)

49th JANNAF Propulsion Meeting (Tucson, AZ, 14-16 Dec 1999)

(Statement A)



BRIEFING FOR INDUSTRY 16 Dec 1999



Basic Research in Supercritical Combustion

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6.1 Objectives

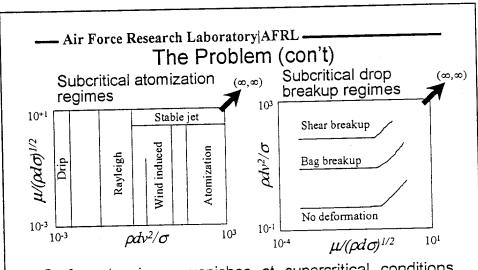
 Determine the mechanisms which control the breakup, transport, mixing, and combustion of suband super-critical droplets, jets, and sprays.

6.1 Funding (\$1,000's) Prior FY99 FY00 FY01 141 141

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The Problem

- It is often advantageous to operate combustion chambers at pressures exceeding the critical pressure of one or both propellants.
 - Higher chamber pressures lead to greater performance (Isp).
- At supercritical pressures, the distinct difference between gas and liquid phases disappears.
 - Conventional "spray combustion" experience no longer applies.
- It is not known how to replace conventional "spray combustion" models in engine design codes.
 - The lack of understanding leads to potentially large engine design errors.



Surface tension σ vanishes at supercritical conditions. Conventional atomization and breakup parameters become *infinite*, where no data exists.

Supercritical atomization and breakup regimes are unknown.

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The Problem (con't)

- Supercritical combustion is complicated by several factors not present in subcritical combustion:
 - Vanishing surface tension.
 - Equivalent gas and liquid phase densities.
 - Strongly enhanced gas / liquid solubility.
 - Different reaction kinetics.
 - Mixing induced critical point variations.
 - Property computation / singular behavior.
 - Zero enthalpy of vaporization.
 - Infinite specific heat (Cp).
 - Infinite compressibility.
- Deeply fundamental questions such as whether droplets can even exist were hotly debated when this work began.

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Technical Approach

- Windowed pressure vessel operating at supercritical pressures.
- Cryogenic fluid capability (LOX, LN2)
- Capability to produce supercritical droplets and jets.
- Shadowgraph, Schlieren, and Raman visualization of concentration fields.
- Capability to drive flows with an acoustic driver

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Payoffs

Provide alternatives to trial and error development

- <u>Performance</u>: Injector related design uncertainties translate to 3-6 sec lsp on a booster class LOX/H2 engine.
 - Comparison: IHPRPT 2010 lsp objective is 13.5 sec.
 - 3-6 sec Isp buys 1.6 3.3 <u>tons</u> payload on the Space Shuttle Main Engine (SSME) worth \$20-40M <u>per launch.</u>
- Operability and Lifetime: Injector related performance deficit required SSME turbopumps to be run at 105% rated power, increasing pump stress.
 - Pumps are the most expensive SSME maintenance item.
 - Turb. blade cracking problem is also probably inj. related.
- <u>Instability</u>: Injector related Saturn F-1 instability problem required over 800 full scale tests to solve.
 - Present day costs: over \$750K per test. Total: \$600 million.

<u>Trial-and-error approaches risk significant cost overruns</u> that can no longer be afforded

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FY99 Accomplishments

For subcritical and supercritical mixing layers:

- Measured the growth rate for a wide variety of propellant combinations
- Developed comprehensive model to predict mixing layer growth rates over four orders of magnitude in density ratio
- Performed fractal analysis of mixing layer geometry
- Installed and performed initial Raman measurements of species distributions.

OUT.

Public

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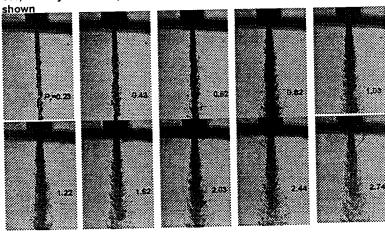
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Evolution of Mixing Layers in Transition from a Subcritical to a Supercritical State

Acknowledgements:
Bruce Chehroudi
Rich Cohn
Roger Woodward
Ed Coy

N₂ into N₂

Back-illuminated images. Chamber is at a fixed supercritical temperature of 300 K but varying sub- to supercritical pressures (P_{critical} = 3.39 MPa). Re = 25,000 to 75,000. Injection velocity: 10-15 m/s. Froud number = 40,000 to 110,000. Injectant temperature = 99 to 120 K. Reduced pressures are



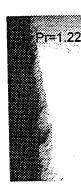


PRESSURE DEPENDENT MIXING LAYER STRUCTURE

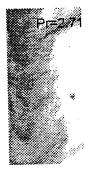
Nitrogen/nitrogen system ($P_{cr} = 3.39 \text{ MPa}$, $T_{cr} = 126 \text{ K}$) T_{inj} = 128 K, T_{amb} = 300 K, mass flow = 350 mg/s



Low Pres. Subcritical Droplets



Mod. Pres. Supercritical Ligaments



High Pres. Supercritical Gas layers

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Sub- and Super-critical Mixing Layer Physics Chehroudi et. al., AIAA 99-0206, AIAA 99-2489

- N2 Jet Into N2 Darkcore (*)
- Steady Diesei-Type Spray L/D=4 ---- Steady Diesei-Type Spray L/D=85
- N2 jet into N2 L/D=200 (*)

- ◆ Cold He jet into N2; L/D=200 (*)
- Cold N2 jet into He; L/D=200 (*)

- Δ O2 jet into N2; L/D=200 (*)
- O2 jet into N2; Darkcore (*)

- Theory (Papamoschou&Roshko)
 - 1.00 Spray Gas/gas regime Transition 0.10 (θ Supercritical Subcritical 0.01 10.00 1.00 0.01 Chamber/Injectant Density Ratio

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Characteristic Times

- Characteristic bulge formation time (τ_b) at the jet interface (Tseng et al.): $(\rho_l L^3/\sigma)^{1/2}$; ρ_l , L, σ are liquid density, characteristic dimension of turbulent eddy, and surface tension, respectively.
- Characteristic time for gasification (τ_g) (D-square law): D²/K; D and K are drop diameter and vaporization constant.
- A Hypothesis: If these two characteristic times (calculated for appropriate length scales) are comparable then an interface bulge may not be separated as an unattached entity (onset of the gas-jet behavior at supercritical condition)

- Theoretical isothermal liquid spray growth rate (θ_s) based on Orr-Sommerfeld equation and stability analysis to find the wavelength of the most unstable interface wave: $\theta_s \cong 0.27 \left[O + (\rho_s/\rho_s)^{0.5} \right]$
- Papamoschou/Rashko theory for incompressible variable-density gaseous mixing layer/jet: $\theta_{P/R} \cong 0.17 \left[1 + (\rho_g/\rho_l)^{0.5}\right]$
- Dimotakis theory for incompressible variable-density gaseous mixing layer/jet: $\theta_{\text{D}} \cong 0.212 \left[0.59 + (\rho_{\text{a}}/\rho_{\text{l}})^{0.5}\right]$
- ALL HAVE THE SQUARE ROOT OF DENSITY RATIO AND THE SAME EQUATION FORMAT

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Correlation

 Based of the information of the previous slide the following "intuitive/smart" equation is proposed for both sub- and supercritical measured growth rates:

$$\theta_{Ch} \cong O.27 \left[\left(\tau_b / (\tau_b + \tau_g) \right) + (\rho_g / \rho_1)^{O.5} \right]$$

Note:

- For isothermal liquid case: $\tau_g >> \tau_b$ and $\tau_g \to \infty$. It then collapses to the isothermal spray case.
- For subcritical the $(\tau_b/(\tau_b+\tau_g))$ is calculated until it reaches 0.5. After that it is maintained constant at 0.5 for supercritical gas-like jet. The transition point is found to be approximately when $(\tau_b/(\tau_b+\tau_g))\cong 0.5$ (i.e. $\tau_b\cong \tau_g$).

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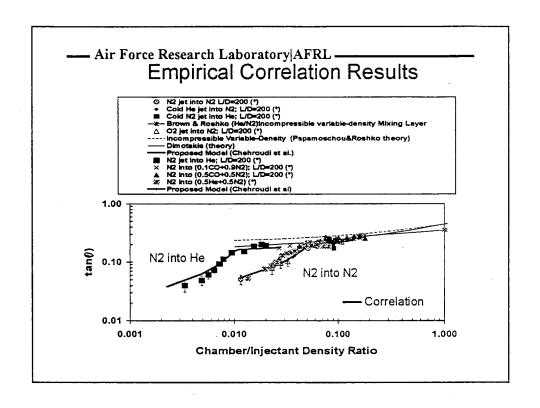
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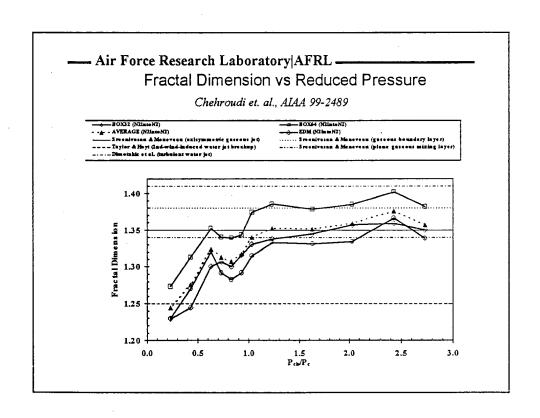
- $(\tau_b/(\tau_b + \tau_a))$ is assumed to be a dominant function of the density ratio (ρ_a/ρ_l) ; i.e. $\tau_b/(\tau_b + \tau_a)) = F(\rho_g/\rho_l)$.
- The function F is only calculated for the N2-into-N2 case and is taken to be the same for other (N_2 -into-He ans N_2 -into-Ar) cases. That is, for example, for N_2 -into-He:

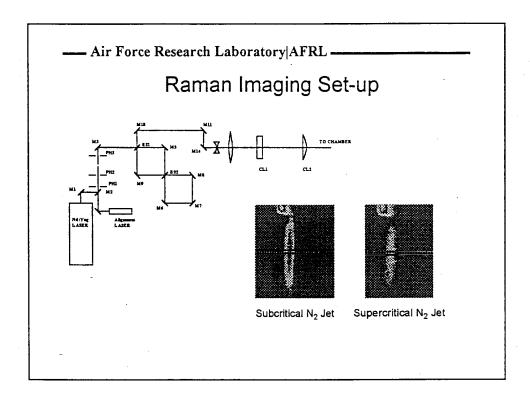
$$\theta_{Ch} {\cong 0.27} \; [\; G(\rho_g/\rho_l) + (\rho_g/\rho_l)^{0.5}] \; \; \text{where} \quad G(\rho_R) = F(\rho_R) \; . \label{eq:theta_chi}$$

$$\rho_R = (\rho_g/\rho_l)$$
; $\rho_R' = \rho_R - (1-X)\rho_R = X\rho_R$

X=1.0 for N_2 -into- N_2 ; X=0.2 for N_2 -into-He; X=1.2 for N_2 -into-Ar.



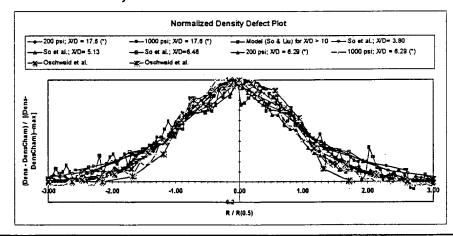




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Preliminary Raman Results

- Plot contains a theoretical model, supercritical jets from AFRL and DLR, and gas jets
- Self-similarity behavior is observed



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Planned for FY00

- Complete Raman species measurements; reduce and analyze data.
- Install acoustic drivers and investigate the effect of acoustic waves.

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Summary and Conclusions

- Structural differences in cryogenic jets have been observed below and above the thermodynamic critical point.
- Liquid-Jet like appearance occurs up to near the critical point, similar to second wind-induced liquid jet breakup regime.
- Gas-jet like appearance occurs above the critical point. No drops are observed.
 - Supercritical spreading rate measurements agree quantitatively with incompressible variable density mixing layer experiments and theory.
 - Supercritical fractal dimensions agree quantitatively with gas jet measurements.
- New and existing mixing layer growth rate experiments and theory have for the first time been consolidated into a single plot as a function of density ratio, where the density ratio spans three orders of magnitude.
- A physical mechanism and correlation have been proposed to describe the transition from spray to gas jet behavior.

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Summary and Conclusions (con't)

 Preliminary analysis of Raman data indicates selfsimilar spreading behavior much like a gas jet.